



Raster Image Correlation Spectroscopy RICS

Novel Idea: Raster Image Correlation Spectroscopy

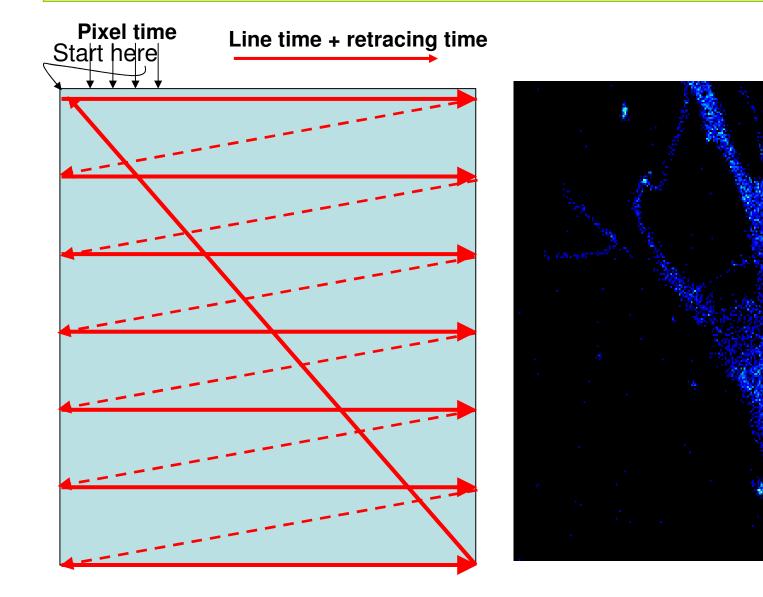
We can have a combination of very high time resolution with sufficient spatial resolution.

Major benefits of RICS:

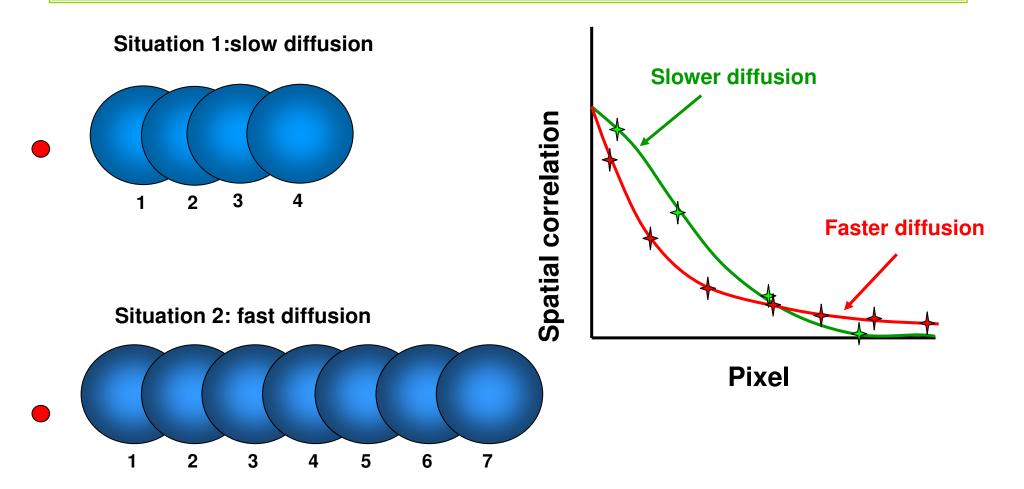
- It can be done with commercial laser scanning microscopes (either one or two photon systems)
- It can be done with analog detection, as well as with photon counting systems, although the characteristic of the detector must be accounted for (time correlations at very short times due to the analog filter)
- > RICS provides an intrinsic method to separate the immobile fraction
- > It provides a powerful method to distinguish diffusion from binding

How does it work?

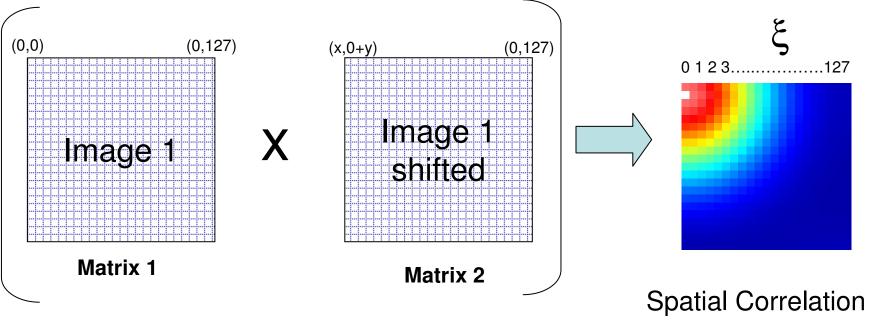
Raster Scanning



Temporal information hidden in the raster-scan image: the RICS approach



How is the spatial correlation done?



Operation:

Results

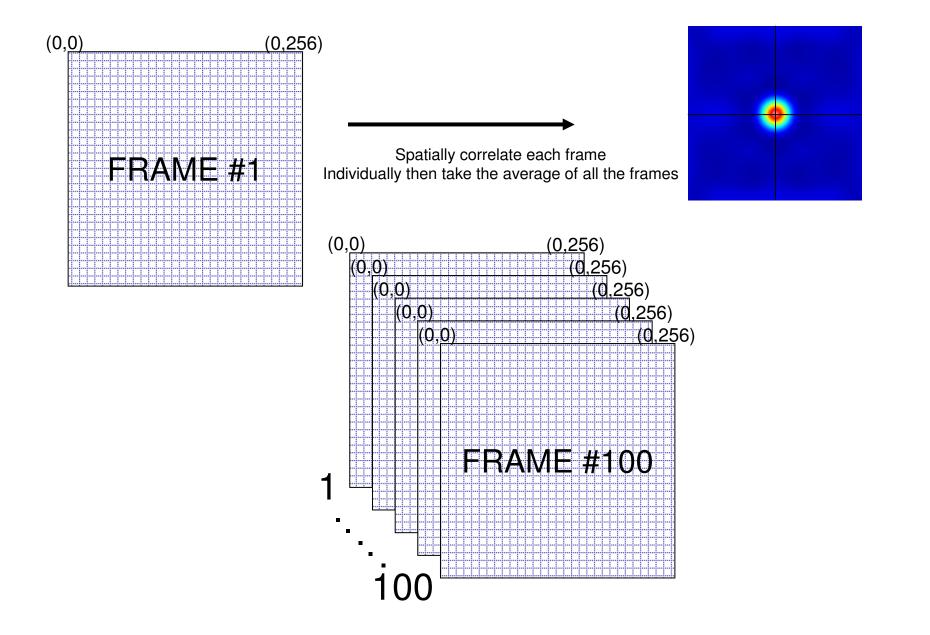
In the x direction

PLUS In the y direction $(0,0\times0,0) + (0,1\times0,1) + (0,2\times0,2)...(0,127\times0,127)$

 $+(1,0\times1,0)+(1,1\times1,1)+(1,2\times1,2)...(1,127\times1,127)$

One number is obtained for x and y and is divided by the average intensity squared

How to use a stack of images?



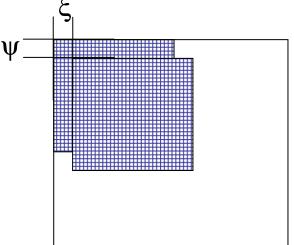
The RICS approach: 2-D spatial correlations

In a raster-scan image, points are measured at different positions and at different times simultaneously

If we consider the time sequence, it is not continuous in time If we consider the pixel sequence, it is contiguous in space

In the RICS approach we calculate the 2-D spatial correlation function (similarly to the ICS method of Petersen and Wiseman)

$$G_{RICS}(\xi, \psi) = \frac{\langle I(x, y)I(x + \xi, y + \psi) \rangle}{\langle I(x, y) \rangle^2}$$



The variables x and y represent spatial increments in the x and y directions, respectively

2-D spatial correlation can be computed very efficiently using FFT methods.

To introduce the "RICS concept" we must account for the relationship between time and position of the scanning laser beam.

The RICS approach for diffusion

The the dynamics at a point is independent on the scanning motion of the laser beam

$$G_{RICS}(\xi, \psi) = S(\xi, \psi) \times G(\xi, \psi)$$

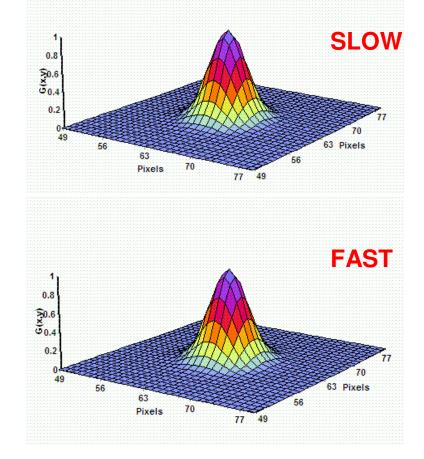
Consider now the process of diffusion. The diffusion kernel can be described by the following expression

$$P(r,t) = \frac{1}{(4\pi Dt)^{3/2}} \exp(-\frac{r^2}{4Dt})$$

There are two parts:

- (1) the temporal term,
- (2) the spatial Gaussian term

For any diffusion the amplitude decreases as a function of time and the width of the Gaussian increases as a function of time



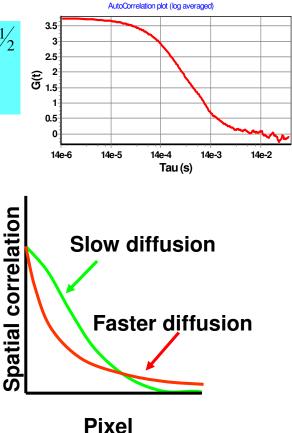
RICS: space and time relationships

At any position, the ACF due to diffusion takes the familiar form:

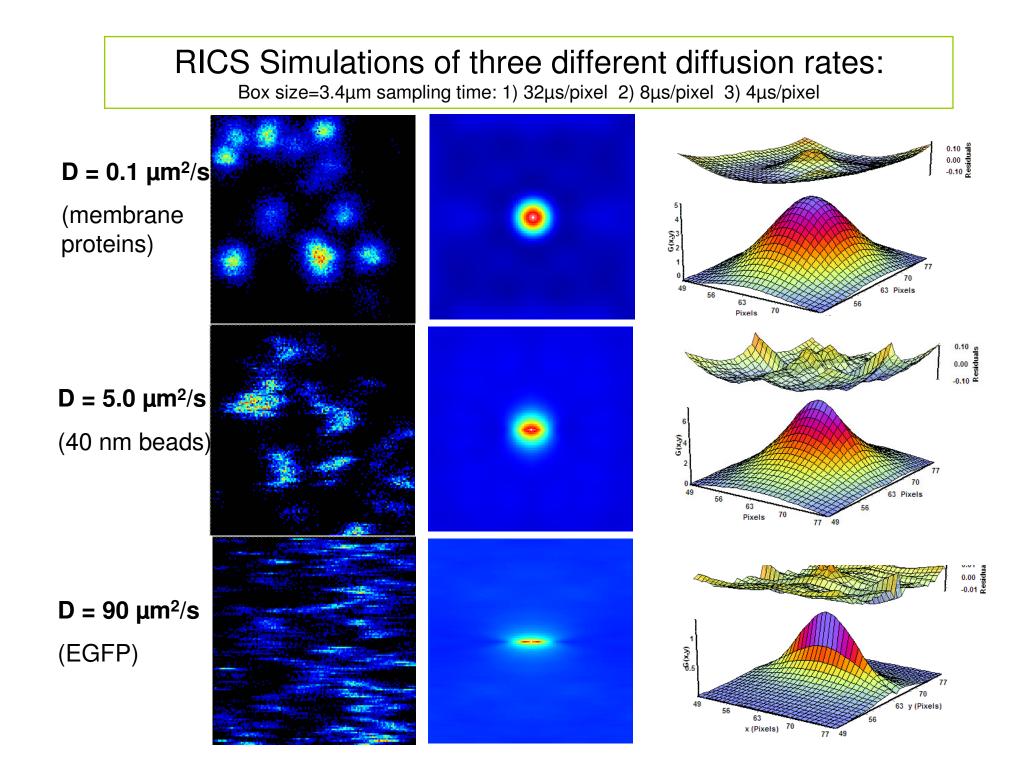
$$G(\xi, \psi) = \frac{\gamma}{N} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_0^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1 + \frac{4D(\tau_p \xi + \tau_l \psi)}{w_z^2} \right)^{-1} \left(1$$

 t_p and t_l indicate the pixel time and the line time. The correlation due to the scanner movement is:

$$S(\xi, \psi) = \exp\left(-\frac{\left[\left(\frac{2\xi\delta r}{w_0}\right)^2 + \left(\frac{2\psi\delta r}{w_0}\right)^2\right]}{\left(1 + \frac{4D(\tau_p\xi + \tau_l\psi)}{w_0^2}\right)}\right)$$



Where δr is the pixel size. For D=0 the spatial correlation gives the autocorrelation of the PSF, with an amplitude equal to γ/N . As D increases, the correlation (G term) becomes narrower and the width of the S term increases.

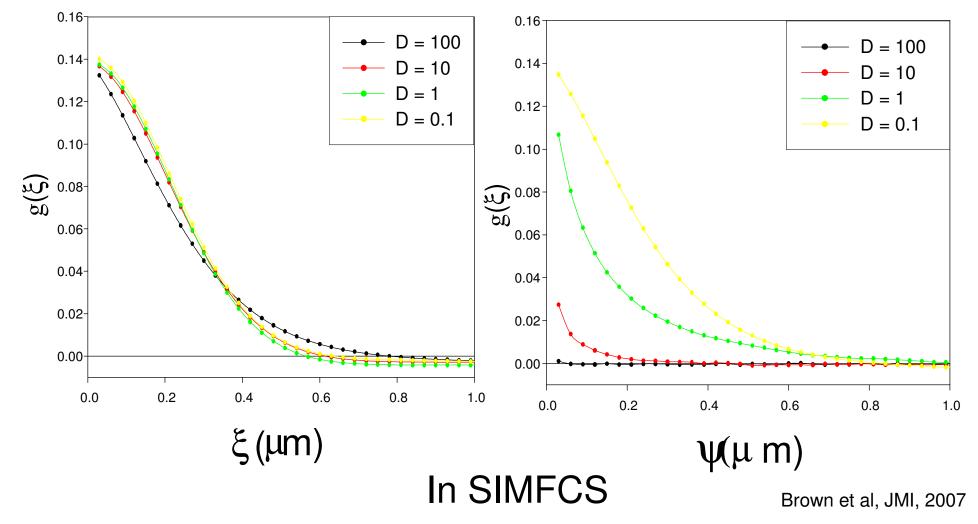


Horizontal and Vertical fits:

Simulations of beads 300 frames, 128x128pixels, 8µs/pix, size of pixels=30nm

Horizontal ACF





How to Setup the Laser Scanning Confocal Microscope

➤ Scan Speeds (µs/pixel):

- $4\mu s$ for fast molecules $D > 100\mu m^2/s$
- 8 -32 μ s for slower molecules D= 1 μ m²/s-100 μ m²/s
- 32-100 μ s for slower molecules D= 0.1 μ m²/s-10 μ m²/s

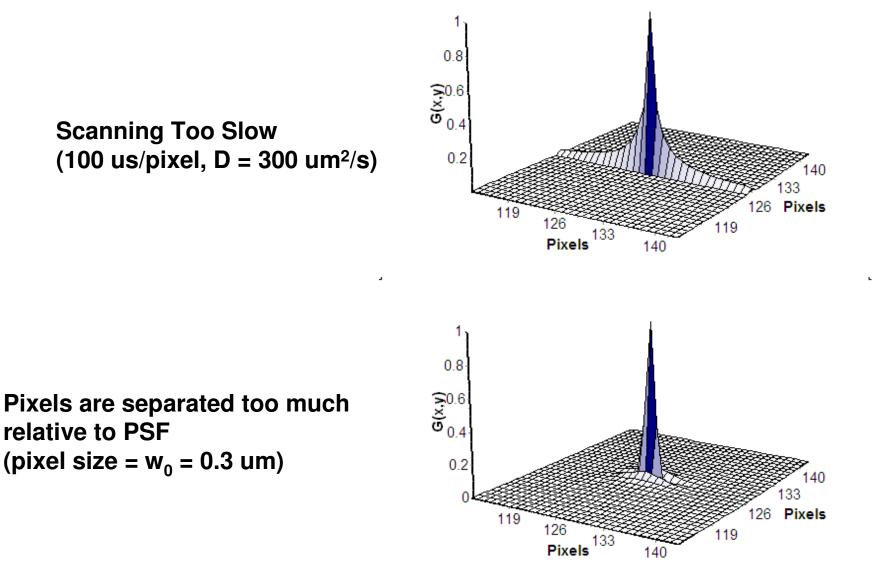
Pixel Size:

• 3-4x smaller than the Point Spread Function (PSF > 300nm)

Molecular Concentrations

•Same conditions as conventional FCS methods

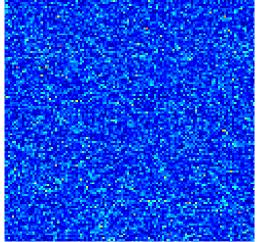
Common Errors in RICS



Courtesy of Jay Unruh

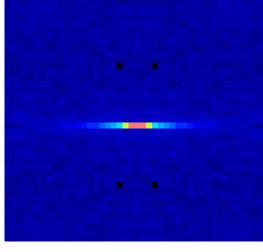
RICS: Fits to spatial correlation functions Olympus Fluoview300 LSM

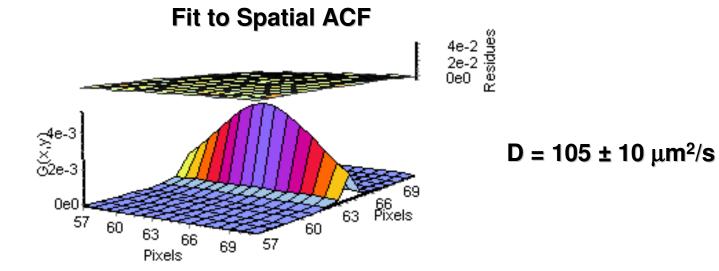
EGFP in solution



128x128, 4 µs/pixel, 5.4 ms/line, 0.023 µm/pixel

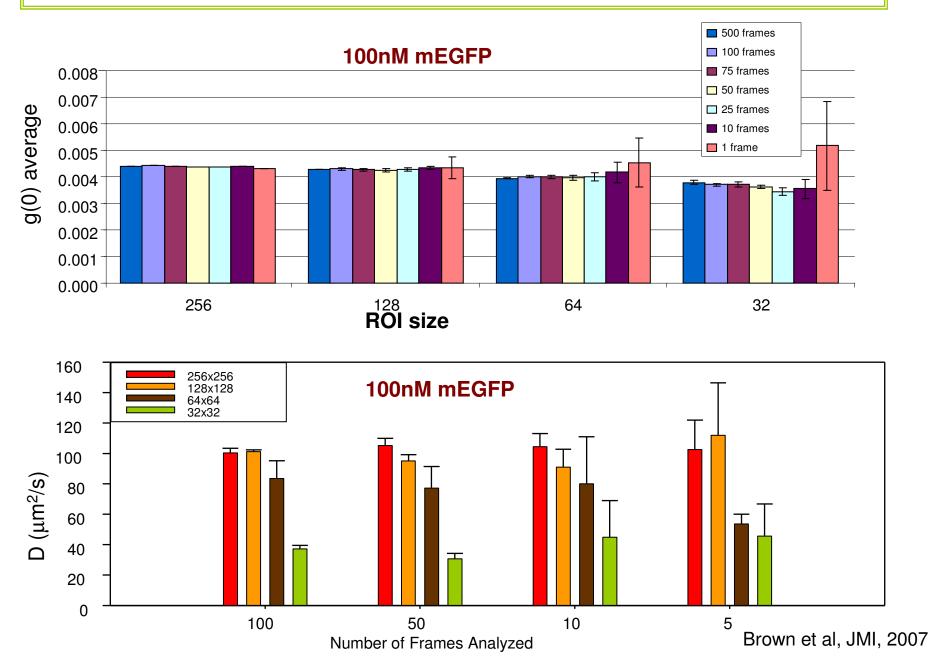




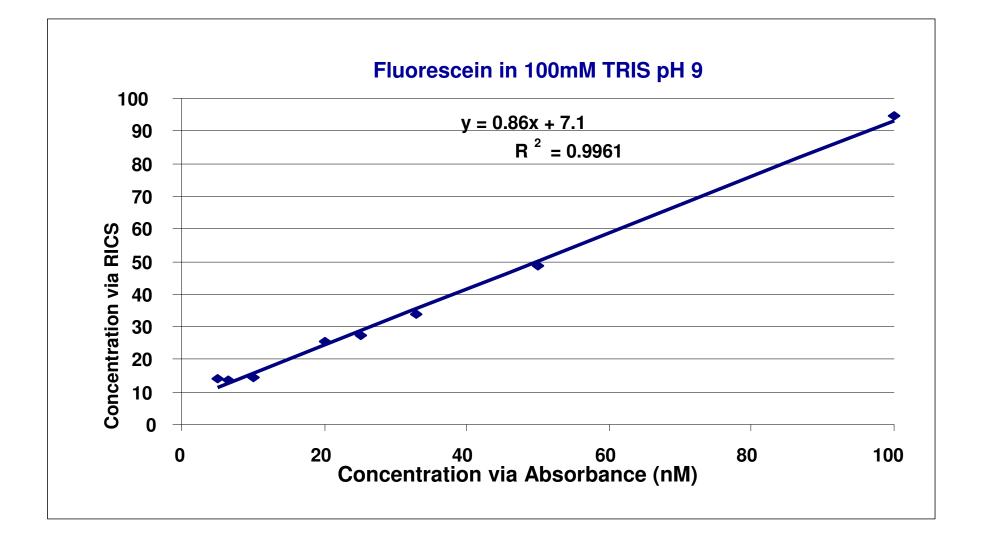


Digman et al. Biophys. J., 2005

What ROI size to use? How many frames to acquire?



Obtaining concentration from RICS



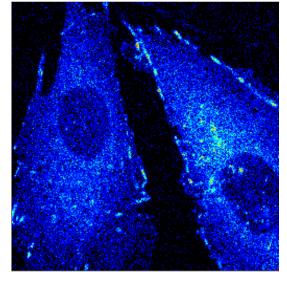
Brown et al, JMI, 2007

How we go from solutions to cells?

In cells we have an immobile fraction

The 2-D-spatial correlation of an image containing immobile features has a very strong correlation pattern

We need to separate this **immobile** fraction from the mobile part before calculating the transform



How is this achieved?

Does noise from the detectors correlate?

In a "truly immobile" bright region, the intensity fluctuates according to the Poisson distribution due to shot noise.

The time correlation of the shot noise is zero, except at time zero.

The spatial correlation of the intensity at any two pixels due to shot noise is zero, even if the two points are within the PSF.

If we subtract the average intensity and disregard the zero time-space point, the immobile bright region totally disappear from the correlation function

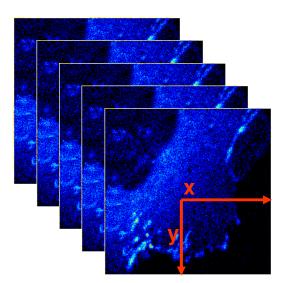
Attention!!!!

This is not true for analog detection, not even in the first order approximation. For analog detection the shot noise is time (and space) correlated. Photon counting: ACF of a bright immobile particle

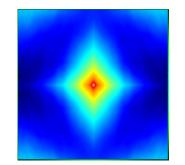


Analog detection: ACF of a bright immobile particle

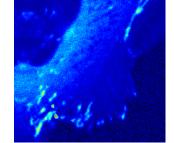
Formula used to subtract background:



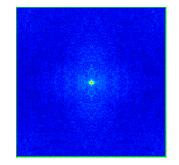
Spatial Correlation



Spatial correlation before subtracting background



Subtract the average



Spatial Correlation of entire image After subtracting image

Average intensity of each pixel on the overall stack: I(x, y)

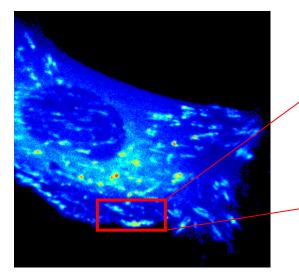
 $I_i(x, y) - \overline{I(x, y)}$

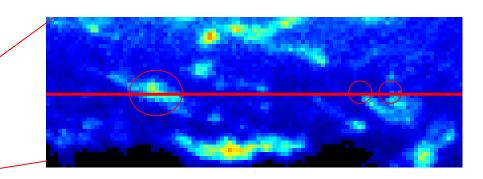
The intensity of each pixel minus the average intensity from entire stack for each pixel: However, this yields negative values.

A scalar must be added : $a = \overline{I}$

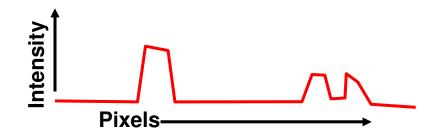
$$ICS(F_i(x, y))$$
 where $F_i(x, y) = I_i(x, y) - \overline{I(x, y)} + a$

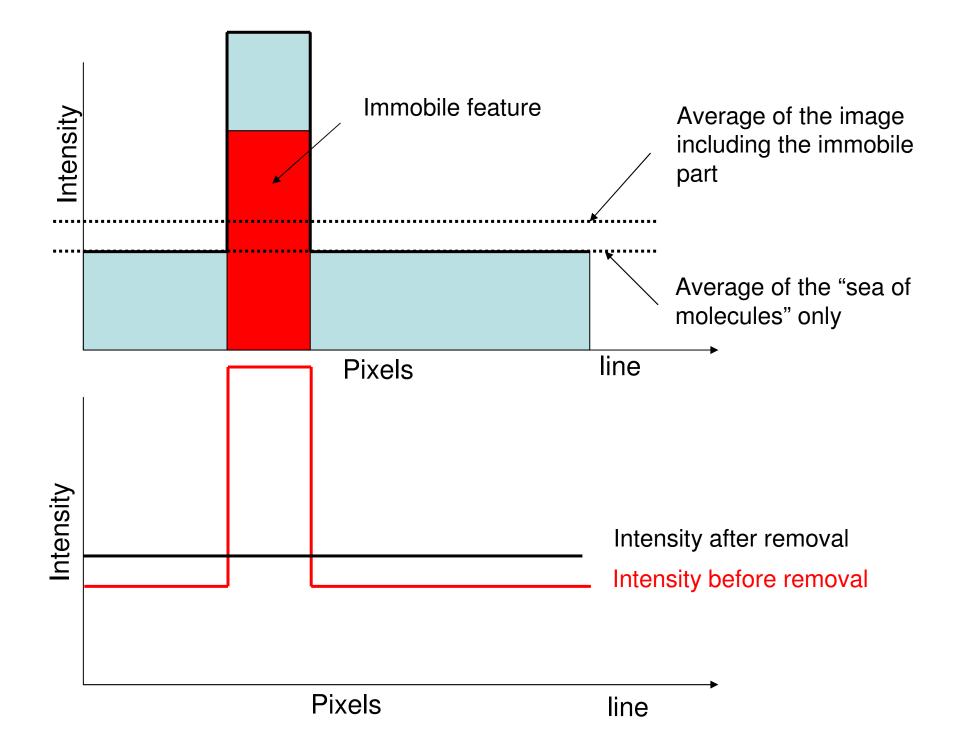
How to subtract immobile features from images?



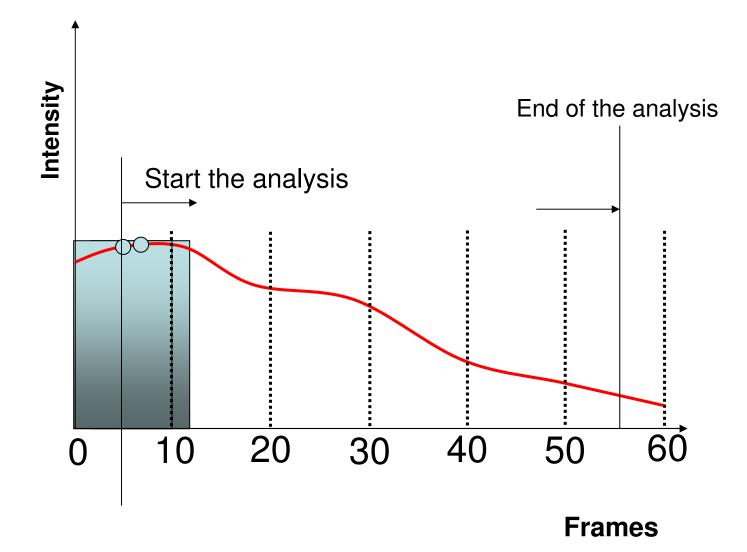


Intensity profile

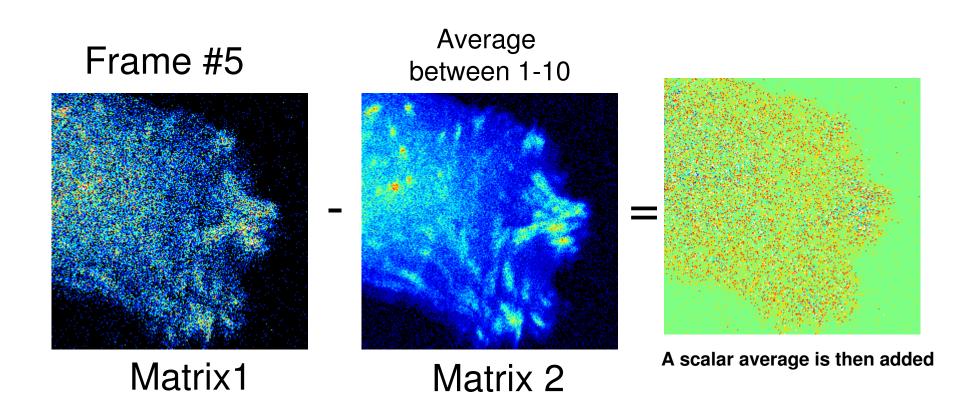




Subtraction of moving average

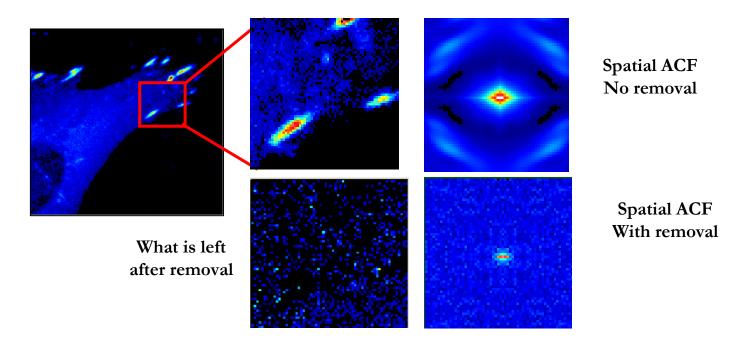


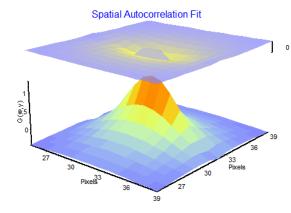
Moving average operation on frames:



Operation is repeated for frame #6 - average between 2-11 frame #7 - average between 3-12

Example of the Removal of Immobile Structures and Slow Moving Features

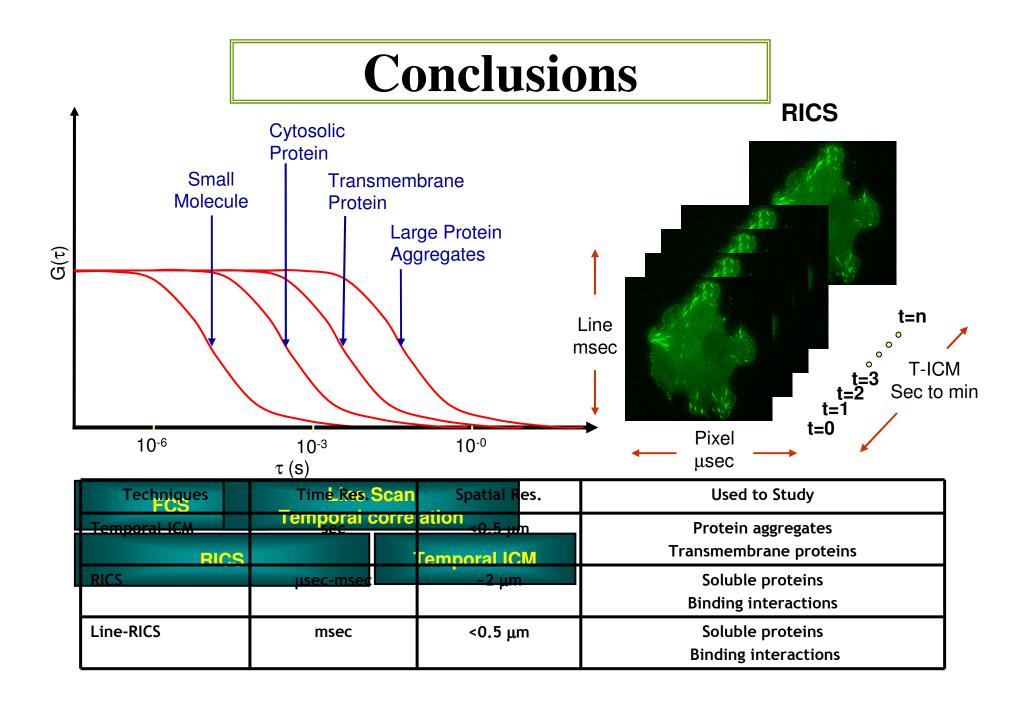




Fit using 3-D diffusion formula

Pixel size =	0.092 µ m
Pixel time=	8 µ s
Line time =	3.152 ms
Wo =	0.35 µ m

G1(0)	=	0.0062	
D1		7.4 µm²/s	
G2(0)	=	0.00023	
D2		0.54 µm²/s	
Bkgd	=	-0.00115	



Summary

- Measures dynamic rates from the µsec-msec time scale
- Anyone with a commercially available instrument can use it
- Immobile structures can be filtered out and fast fluctuations can be detected
- RICS has high spatial and temporal resolution
- The range of these dynamic rates covers a wide range from immobile to cytosolic diffusions (0.2-12um2/s)
- Other types of processes and interactions are also measured
- Line scanning is essential for determination of binding process and complements the RICS analysis

We have expanded the RICS methods to do Cross-Correlation RICS (ccRICS)

The ccRICS approach

The spatial correlation function

$$G_{ccRICS}(\xi, \psi) = \frac{\langle I_1(x, y)I_2(x + \xi, y + \psi) \rangle}{\langle I_1(x, y) \rangle \langle I_2(x, y) \rangle} - 1$$

The variables ξ and ψ represent spatial increments $% \xi$ in the x and y directions, respectively

The $G_{cc}(0,0)$ value and bleedthrough $G_{cc}(0,0) \propto \left[\frac{f_{11} f_{12} \langle N_1 \rangle + f_{21} f_{22} \langle N_2 \rangle}{f_{11} f_{12} \langle N_1 \rangle^2 + (f_{11} f_{22} + f_{21} f_{12}) \langle N_1 \rangle \langle N_2 \rangle + f_{21} f_{22} \langle N_2 \rangle^2} \right]$ 100 80 60 %T 40 Ch1 Ch2 20 $F_1(t) = f_{11}N_1 + f_{21}N_2$ $F_2(t) = f_{12}N_1 + f_{22}N_2$ 0 450 500 550 600 650 700 Wavelength (nm)

Experimental issues

- The volume of excitation and emission at the two excitation wavelengths must superimpose (we are using the Olympus FV1000 LSCM for these experiments)
- Bleedthrough of the green into the red channel must be small (<5%)
- FRET will strongly decrease the ccRICS signal
- High ratio of labeled to unlabeled molecules are needed (if you have only 10% labeled, in a complex of 1:1, you will only have 1% of the complexes labeled with both proteins)

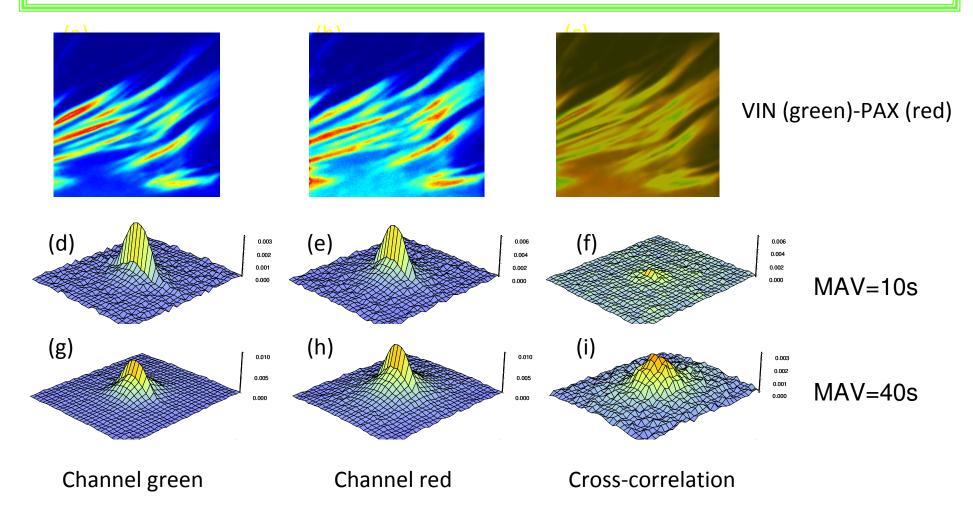
Cells. MEF transfected Vinculin, FAK and paxillin. cDNA were ligated to EGFP or mCherry at the C-terminal end.

Microscopy. Olympus FV1000 with 60x 1.2NA water objective, 12.5 us/pixel, 256x256 pixels 12.5 µm square, 100 to 200 frames collected for each sample. 1frame/s.

EGFP excitation at 488nm (0.5%) and mCherry at 559nm (adjusted to a max of 1.5%). Emission filters at 505-540nm and 575-675 nm, for the green and red channels, respectively.

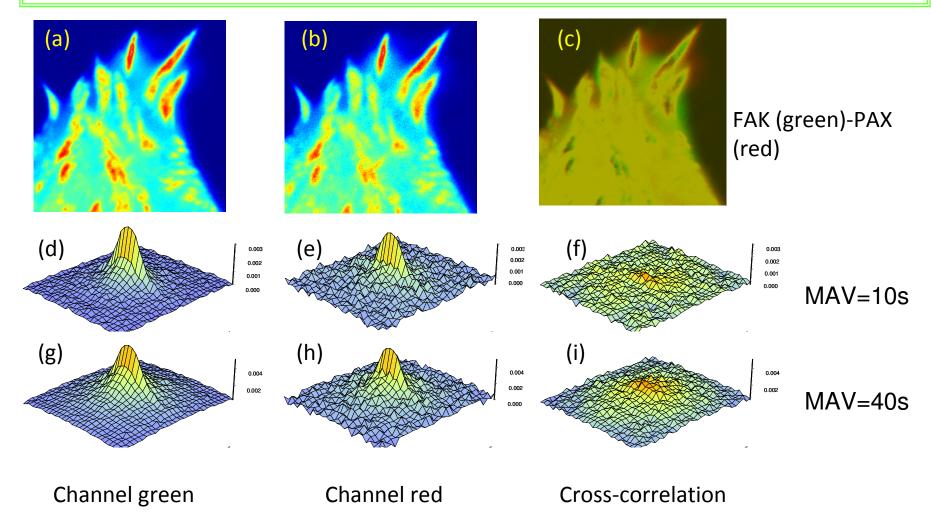
Overlap of the volume of observation was tested by imaging single 100 nm fluorescent beads carrying two colors simultaneously

VIN and PAX co localize at adhesions but they are moving independently in the cytoplasm



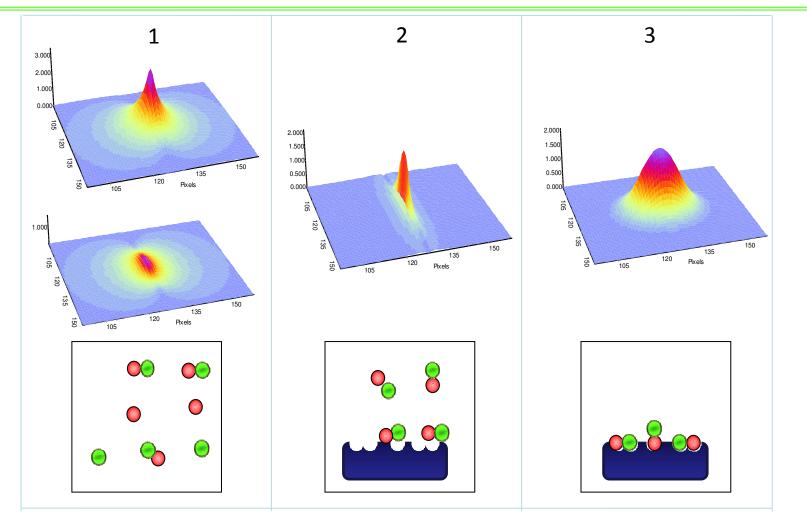
The cross-correlation increases for the slow fluctuations (at MAV=40s). It is round in shape indicating that it is generated at single locations.

FAK and PAX co localize at adhesions but they are moving independently in the cytoplasm



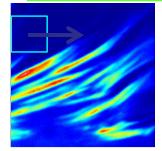
The cross-correlation increases for the slow fluctuations (at MAV=40s). It is round in shape indicating that it is generated at single locations and it is very small.

Schematic representation for the interpretation of the ccRICS experiment. Simulation of binding and diffusion

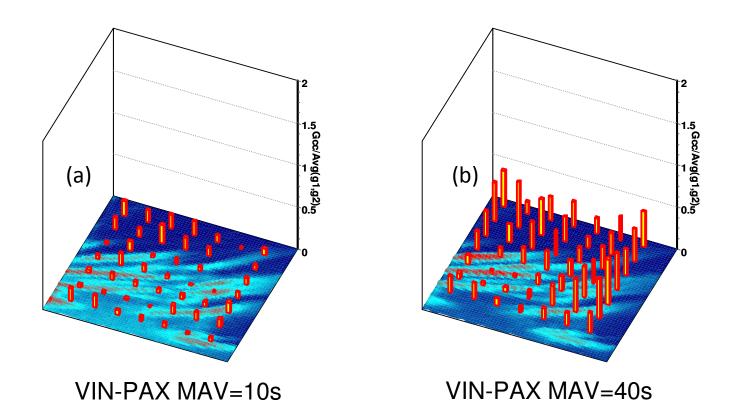


Diffusion Few complexes Fast binding Different shape Smaller than PSF Slow binding Round shape

Distribution of fraction of cross-correlation in the cell. Correlation with adhesion disassembling



ccRICS by scanning a region of interest across the image Calculating the ratio $G_{cc}/AV(G_1,G_2)$

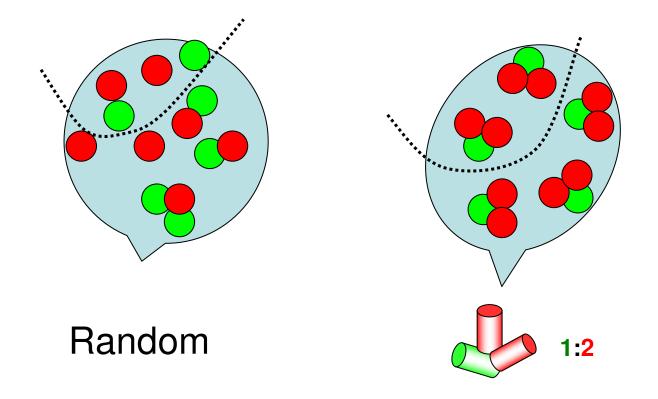


There is "more" cross-correlation at the locations of adhesion disassembling

Summary of ccRICS

- We developed a **toolbox** for biophysicists and cell biologists to address common questions regarding the formation of protein complex, their spatial distribution and their stoichiometry
- ccRICS is extremely powerful at detecting joint diffusing proteins and in separating diffusion from binding processes
- The Paxillin, vinculin and FAK never crosscorrelate in the cytoplasm before binding to the focal adhesion. We only detect cross correlation due to dissociation of large clusters of proteins.

What is the stoichiometry of these clusters and is this stoichiometry crucial for the biological system?



The Number & Molecular Brightness (N&B) Method





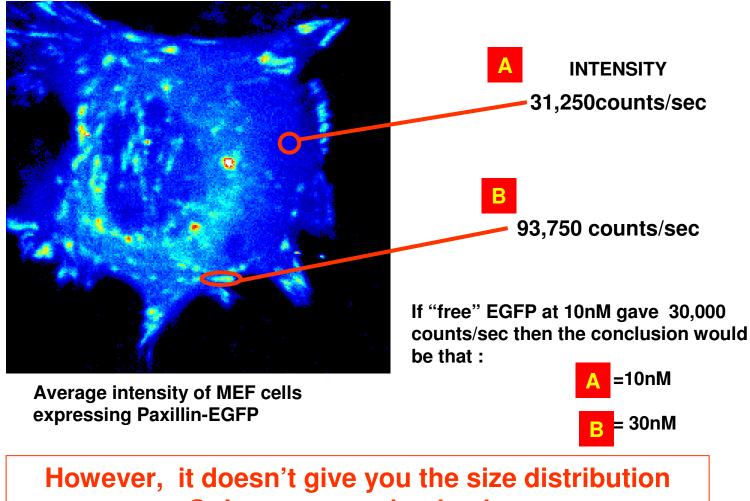






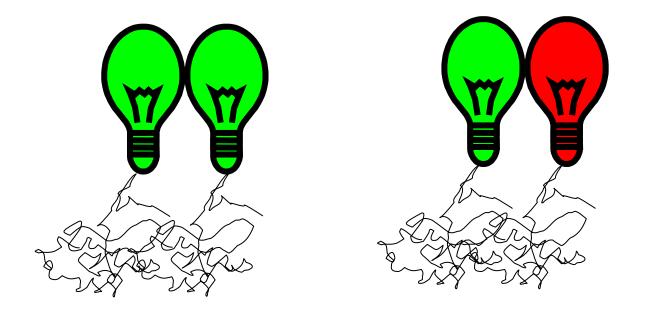
Existing Methods to determine protein concentration and aggregation of proteins in cells

1. Calibration of the free fluorophore based on intensity



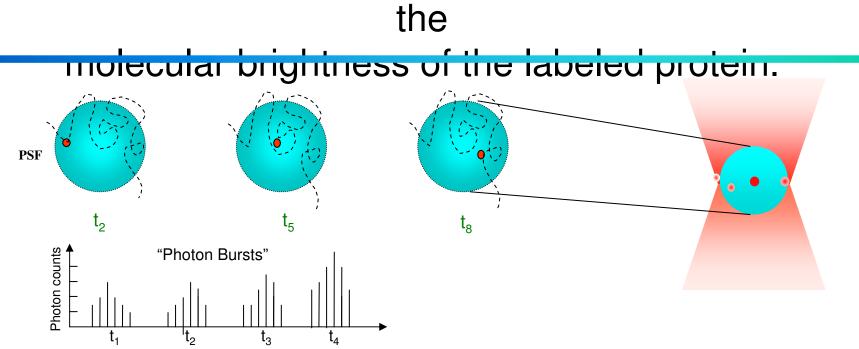
Only concentration is given

2. Förster resonance energy transfer (FRET)



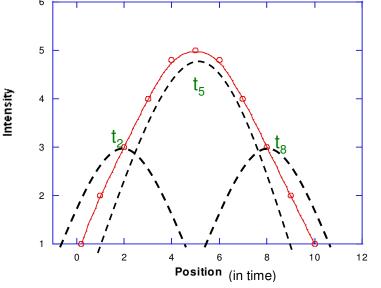
This method is very sensitive to detect the formation of pairs.

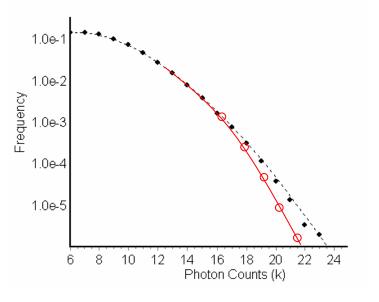
The average photon countrate or pursts determines



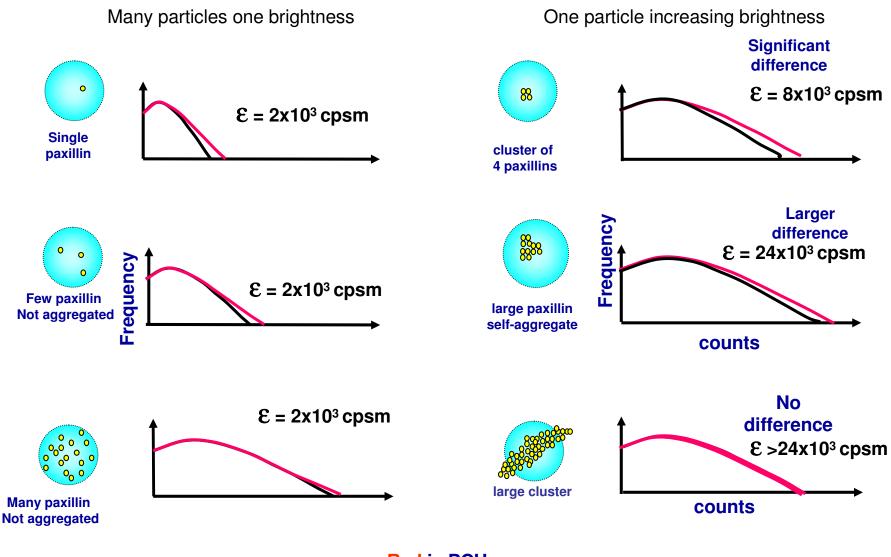
The intensity distribution accounts for the fluctuations of photons from the molecule freely diffusing through the exitation profile. Thus, the overall photon counting count distribution is the weighted superposition of individual Poissonian distributions for each intensity values with a scaling amplitude. The fluctuations light intensity results in a broadeing of photon count distribution with respect to a pure poisson distribution. As the

fluctuations increases, the photon count distribution broadens.



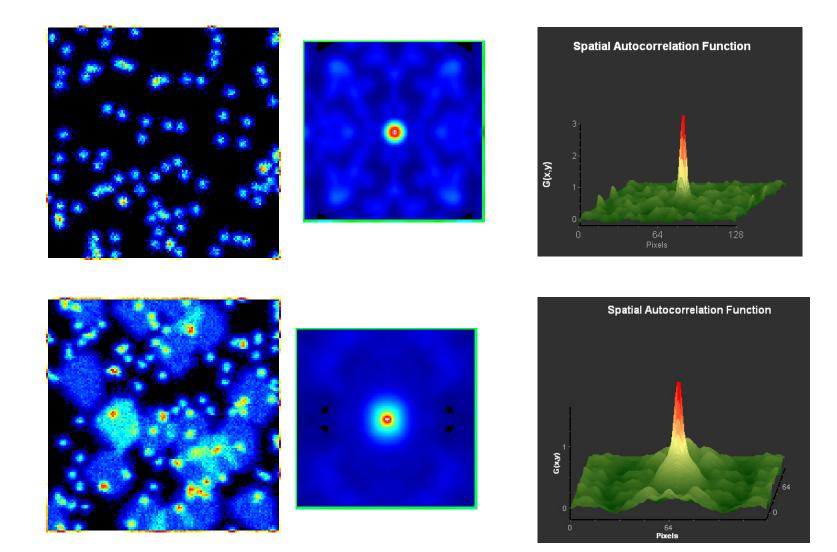


4. Photon Counting Histogram Analysis



Red is PCH Black is Poisson

3. Image correlation Spectroscopy (ICS)



However, the events must be slow >1sec (no movement during one frame) and the aggregates must be large. Petersen and Wiseman:Biophys J. 1999

The Number and Brightness (N&B) analysis

Purpose: Provide a pixel resolution map of molecular number and aggregation in cells

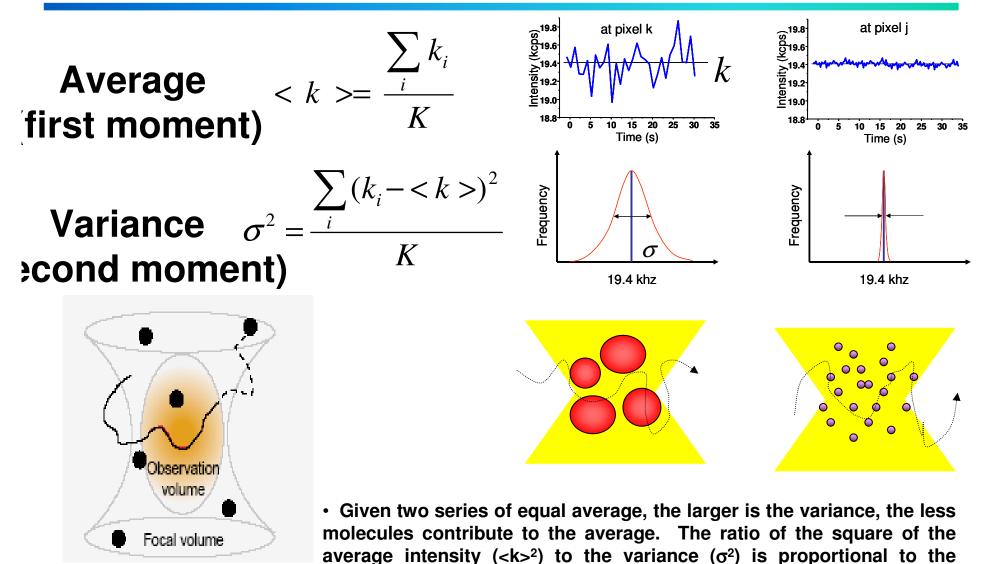
- Method: First and second moment of the fluorescence intensity distribution at each pixel
- Source: Raster scanned image obtained with laser scanning microscopes TIRF with fast cameras Spinning disk confocal microscope

Output: The N and B maps, B vs intensity 2D histogram

Tools: Cursor selection of pixel with similar brightness Quantitative analysis of center and std dev of the *e* and *n* distribution Tools for calibration of analog detectors

Tutorials: mathematical background, data import, analysis examples (our web site

How to distinguish pixels with many dim molecules from pixels with few bright molecules?



average number of particles <N>.

* Originally developed by Qian and Elson (1990) for solution measurements.

 $G(0) = \sigma^2 / \langle k \rangle^2 = 1/N$

Calculating protein aggregates from images

This analysis provides a map of <N> and brightness (B) for every pixel in the image. The units of brightness are related to the pixel dwell time and they are "counts/dwell time/molecule".

$$\langle k \rangle = \frac{\sum_{i} k_{i}}{K}$$
 $\sigma^{2} = \frac{\sum_{i} (k_{i} - \langle k \rangle)^{2}}{K}$

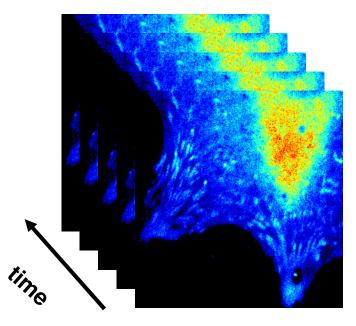
$$B = \frac{\langle k \rangle}{\langle N \rangle} = \frac{\sigma^2}{\langle k \rangle}$$

$$< N >= \frac{< k >^2}{\sigma^2}$$

 σ^2 = Variance

<k>= Average counts

- N = Apparent number of molecules
- B = Apparent molecular brightness
- K = # of frames analyzed



Selecting the dwell time

To increase the apparent brightness we could increase the dwell time, since the brightness is measured in counts/dwell time/molecule.

1 μs dwell time 1 ms dwell time Increasing the dwell time decreases the amplitude of the fluctuation. 0.3 0.2 **G**(1) 0.1 0.0 10^{-5} 1 0⁻³ 10^{-1} 10 10 Tim e(s)

What contributes to the variance?

Variance due to particle number fluctuations

$$\sigma_n^2 = \varepsilon^2 n$$

Variance due to detector shot noise

$$\sigma_d^2 = \mathcal{E}n$$

The measured variance contains two terms, the variance due to the particle number fluctuation and the variance due to $d_{12}^{2} = \sigma^{2} d_{12}^{2} + \sigma^{2} d_{12}^{2}$

S These two terms have different dependence on the $\sigma_n^2 = \varepsilon^2 n$ $\sigma_d^2 = \varepsilon n$ (for the photon counting detector)

Both depend on the intrinsic brightness and the number of molecules. We can invert the equations n is the true number of m and obtain n and ϵ

How to Calculate n and $\boldsymbol{\epsilon}$

$$B = \frac{\sigma^{2}}{\langle k \rangle} = \frac{\sigma_{n}^{2}}{\langle k \rangle} + \frac{\sigma_{d}^{2}}{\langle k \rangle} = \frac{\varepsilon^{2}n}{\varepsilon n} + \frac{\sigma_{d}^{2}}{\langle k \rangle} = \varepsilon + 1$$

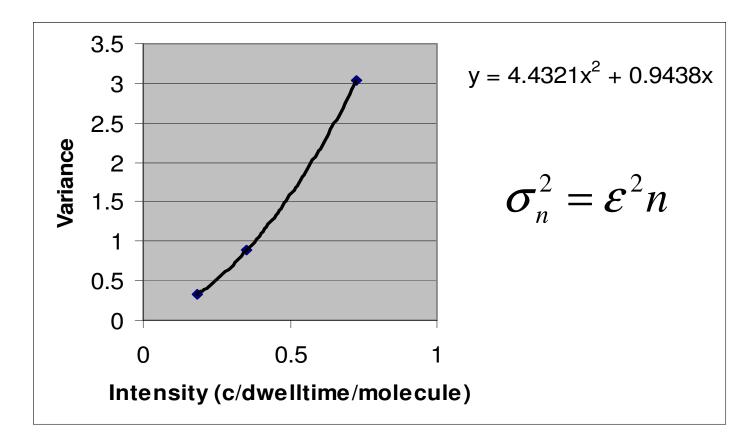
This ratio identifies pixels of different brightness due to mobile particles.

The "true" number of molecules n and the "true" molecular brightness for mobile particles can be obtained from



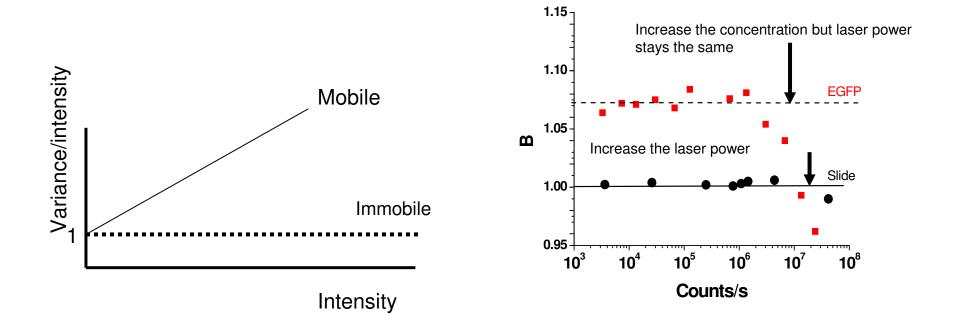
If there are regions of immobile particles, n cannot be calculated because for the immobile fraction the variance is = $\langle k \rangle$. For this reason, several plots are offered to help the operator to identify regions of mobile and immobile particles. Particularly useful is the plot of N*vs*B.

Quadratic dependence of the variance on particle brightness 20nM EGFP in solution as a function of laser power



2-photon excitation using photon counting det

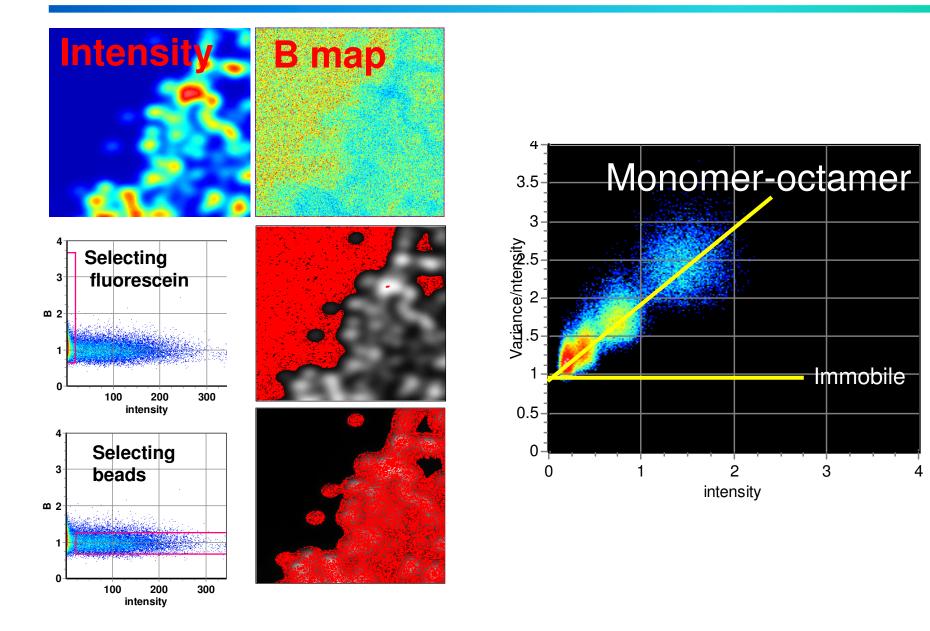
Identification of mobile and immobile molecules



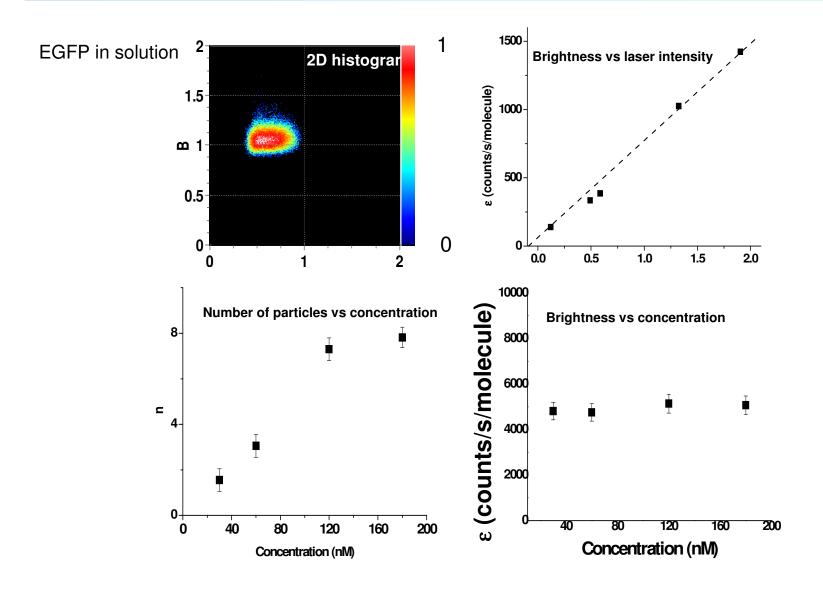
If we change the laser power, a plot of the ratio variance/intensity vs intensity can distinguish the mobile from immobile fraction. The two curves are for different pixel integration times.

The effect of the immobile part: with photon counting detectors

Fluorescent beads in a sea of 100nM Fluorescein.

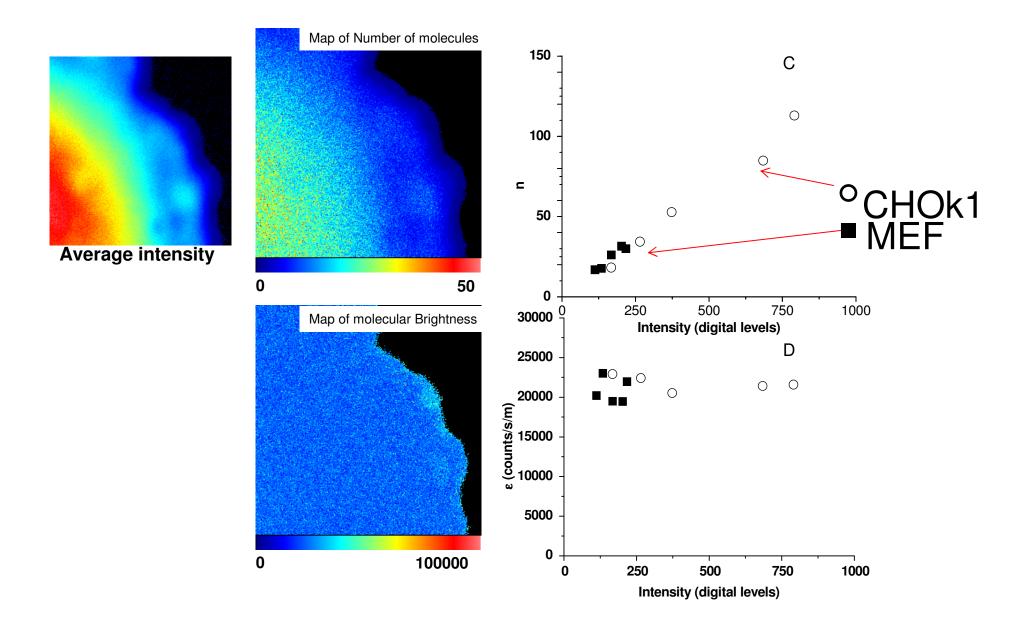


Brightness and number of molecules can be measured independently

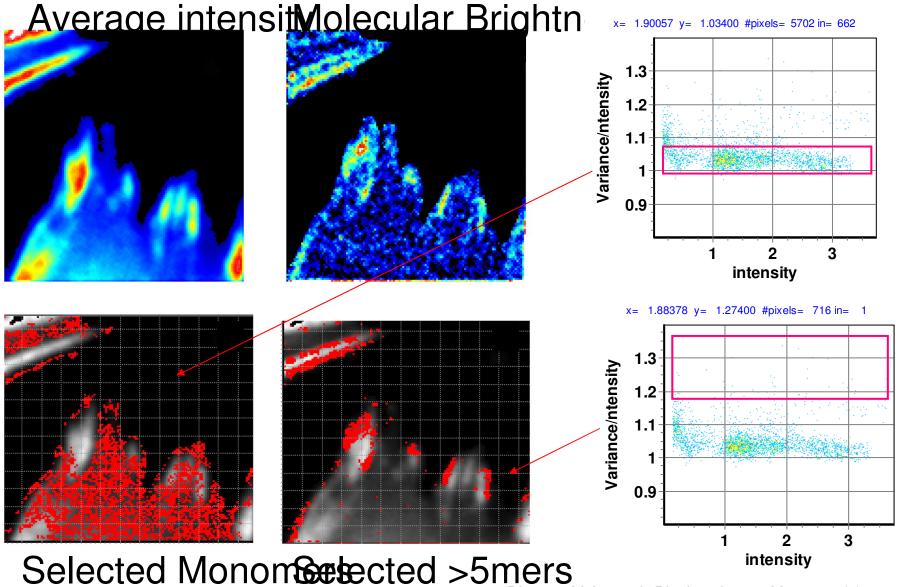


EGFP in CHO-k1 (1-Photon LSM)

homogenous Brightness & heterogeneous Number of Molecules



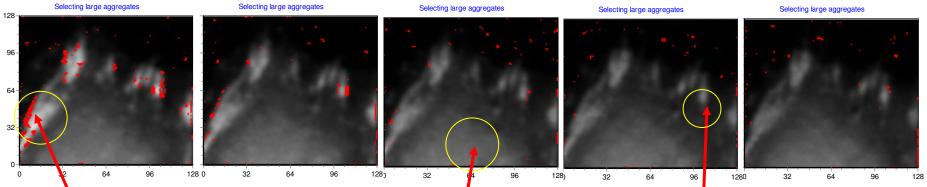
Paxillin assembles as monomers and disassembles as aggregates as large as 8-12



Digman, M.A., et al, *Biophys J.* 2008 Mar 15;94(6):2320-32

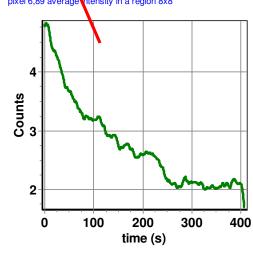
Assembly and disassembly occurs in the order of minutes

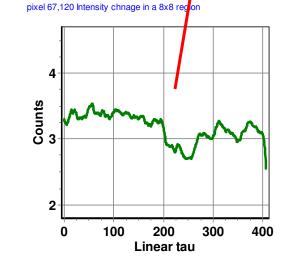
100 Frame average Selecting large aggregates

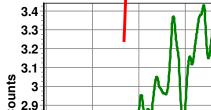


100 to 199200 to 299 300 to 399400 to 499

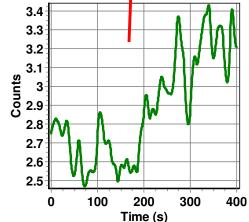








pixel 104,68 A= 0.00000 k= 0.00000 B= 0.00000



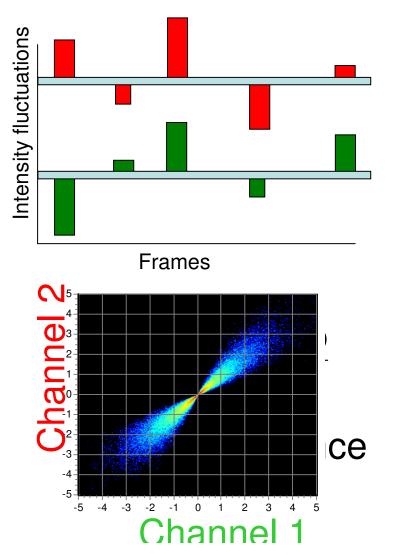
Digman, M.A., et al, *Biophys J.* 2008 Mar 15;94(6):2320-32

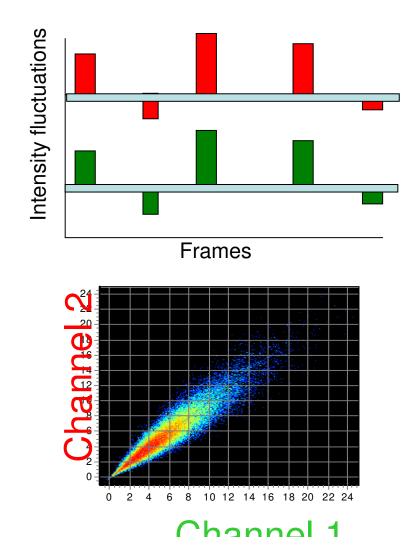
Cross N&B

Conceptual illustration of Cross N&B

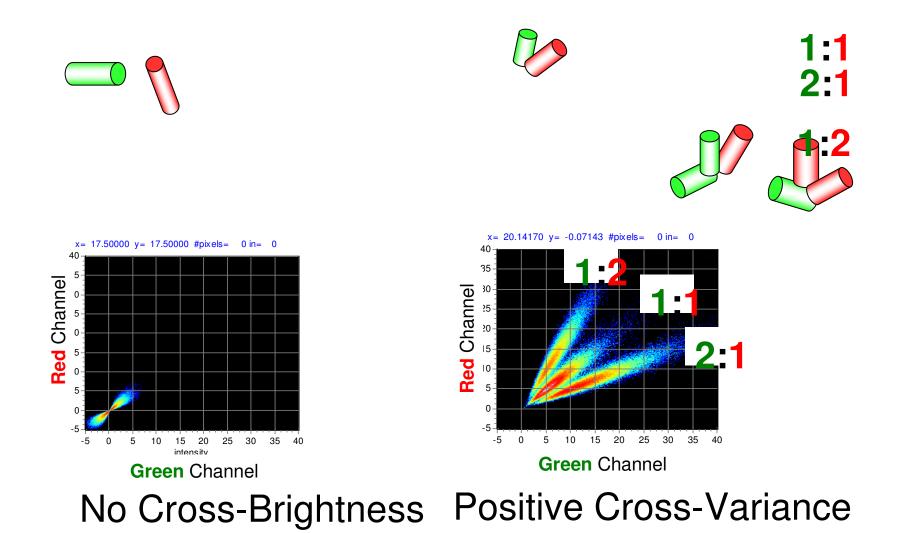
Uncorrelated

Correlated





Cross N&B Analysis determines stoichiometry



This example is only for ideal systems where the brightness is calibrated for both channels.

The co-variance principle and the derivation of the ccN&B method

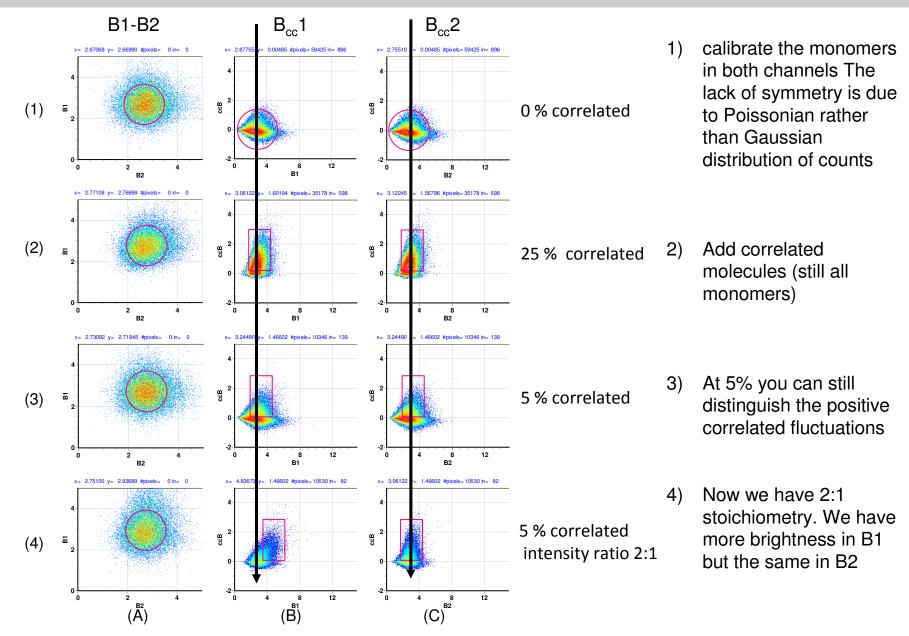
$$\sigma_{cc}^{2} = \frac{\sum (G_{i} - \langle G \rangle)(R_{i} - \langle R \rangle)}{K}$$

$$N_{cc} = \frac{\langle G \rangle \langle R \rangle}{\sigma_{cc}^2}$$

Definition of **co-variance**. It is the average of product of the fluctuations in the Seminion of the **chass number** of molecules. It is the co-variance divided by the product of the intensity in the two channels

K is the number of frames. σ^2 the variance and *<* > indicat

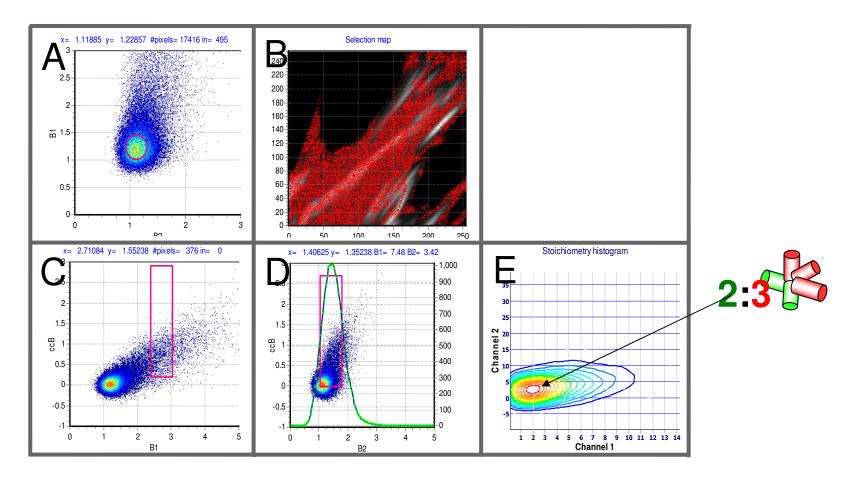
To calibrate the system we need to know the brightness of the monomers



What to look for:

- 1) First we need to calibrate the monomers
- 2) We have to see if there is positive cross variance
- 3) We have to see where the cross variance occurs in respect to the brightness of Ch1 and Ch2

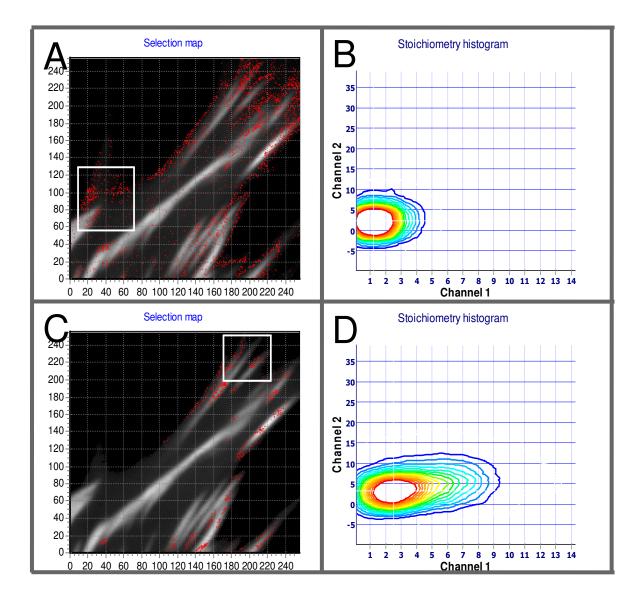
The unknown sample: Vinculin-EGFP and Paxillin-mcherry



Look at the brightnesses that coincide for Ch1 a

We must find for each value of B1 in one pixel, what is the The fluctuations must be correlated so we only look at th

Selecting different regions of the image



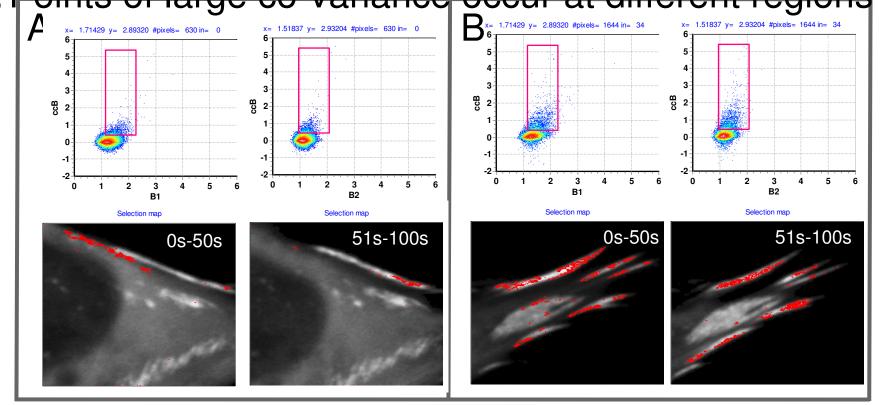


Slide 62

AFH1 Make sure I didnt mess this up. Rick Horwitz, 6/18/2008

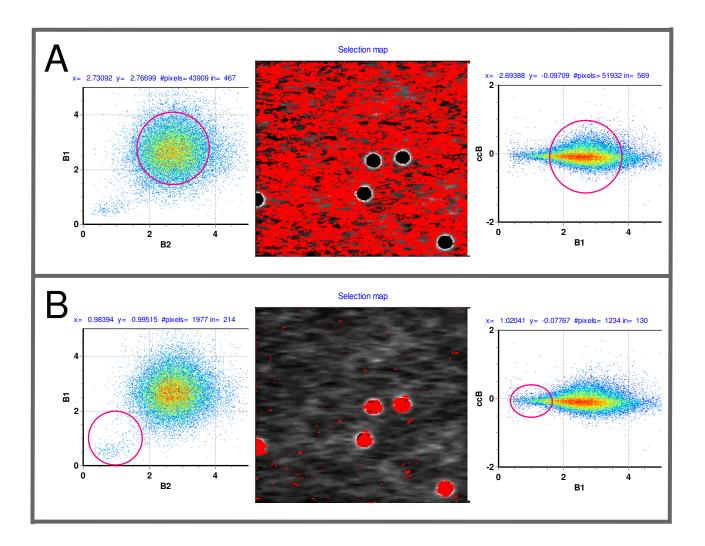
Cross – correlations occur at specific pixels at the adhesions Vinculin-EGFP and Paxillin-mcherry

Large Cross variance is only seen at the adhesion Points of large co-variance occur at different regions an



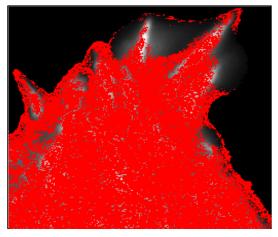
We determined that the brightness for monomers B1 = 1.118 and B2 = 1.22. Thus the ccB1 = 6x monomer and for ccB2 = 3 x the monomer

Simulations: effect of bright immobile features

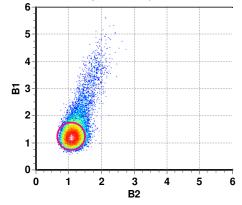


FAK and **Paxillin**

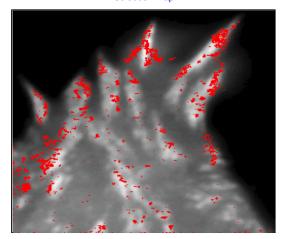
Selection map

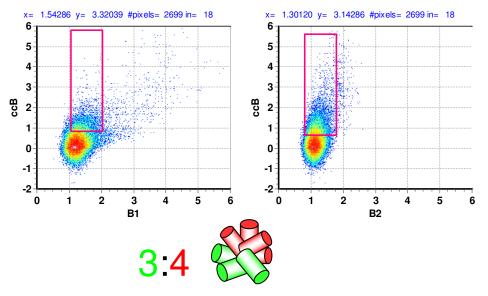


x= 1.08434 y= 1.21053 #pixels=37301



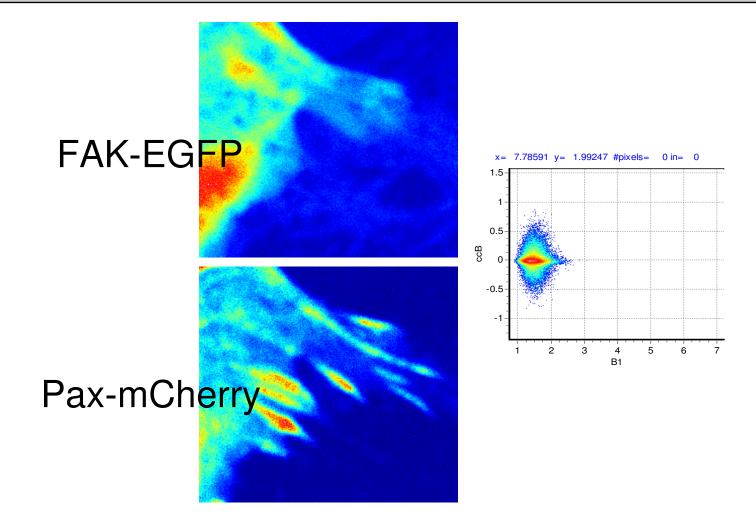
Selection map





Digman, M.A., et al, PNAS Jan.23, 2009 Ahead of print

Testing for artifacts: FAK mutant does not form complexes



mutFAK-PAX cell shows no cross-correlation although the cell

Summary

- N&B distinguishes between number of molecules and molecular brightness in the same pixel
- The acquisition for the N&B can be done with a commercial Laser Scanning Microscope (LSM) and the same data used for RICS can be used to map N and B.
- The Immobile fraction can be separated since it has a Brightness value =1
- The N&B analysis of paxillin at adhesions shows large aggregates of protein during disassembly.
- Cross N&B allows us to determine the stoichiometry of the complexes.

Additional Reading

- 1) Jay R Unruh and Enrico Gratton. Analysis of molecular concentration and brightness from fluorescence fluctuation data with an electron multiplied CCD camera. Biophys J. 2008; [epub ahead of print].
- 2) Michelle A Digman, Rooshin Dalal, Alan R Horwitz, and Enrico Gratton. Mapping the number of molecules and brightness in the laser scanning microscope. Biophys J. 2008; 94(6): 2320-2332.
- 3) Rooshin B Dalal, Michelle A Digman, Alan R Horwitz, Valeria Vetri, and Enrico Gratton. Determination of particle number and brightness using a laser scanning confocal microscope operating in the analog mode. Microsc Res Tech. 2008; 71(1): 69-81.
- Yan Chen, Joachim D Müller, Qiaoqiao Ruan, and Enrico Gratton. Molecular brightness characterization of EGFP in vivo by fluorescence fluctuation spectroscopy. Biophys J. 2002; 82(1): 133-44.
- 5) Alberto Garcia-Marcos, Susana A Sánchez, Pilar Parada, John S Eid, David M Jameson, Miguel Remacha, Enrico Gratton, and Juan P G Ballesta. Yeast ribosomal stalk heterogeneity in vivo shown by two-photon FCS and molecular brightness analysis. Biophys J. 2008; 94(7): 2884-2890.
- 6) Michelle A Digman, Paul W Wiseman, Colin K Choi, Alan R Horwitz, and Enrico Gratton. Mapping the stoichiometry of molecular complexes at adhesions in living cells.Proc Natl Acad Sci USA. 2008; [submitted].